## Problem I

## Reflection and refraction of light

## A. $\mathcal{A n}$ interesting prism

The main section of a glass prism, situated in air ( $n^{\prime}=1.00$ ), has the form of a rhomb with $\Varangle B A D=\Varangle B C D \equiv \theta$. A thin yellow beam of monochromatic light, propagating towards the prism, parallel with the diagonal AC of the rhomb, is incident on the face AB (Fig. 1). The beam is totally reflected on the faces $A D$ and $D C$, then emerges through its face BC. For the yellow radiation, the refraction index of the glass is $n=1.60$.


Fig. 1

| A1 | Derive the mathematical expression for the angle $\theta$ as a function of the <br> refraction index $n$ of the prism, such that the total deviation of the beam that <br> exits the prism to be zero. Under the above condition, calculate the numerical <br> value of $\theta$ in degrees and minutes, if $n=1.60$. | 2.25 p |
| :--- | :--- | :--- |

The prism with $\theta$ determined above and the direction of the incident beam remain fixed, but the nature of the light radiation changes, being formed now of the yellow doublet of the mercury. The two wavelengths have the values 579.1 nm , respectively 577 nm . The refraction indices of the glass for these wavelengths are $n=1.60$, respectively $n+\Delta n$, where $\Delta n=1.3 \cdot 10^{-4}$. The light rays that exit the prism enter longitudinally into an astronomical telescope adjusted for infinite distance.

| A2 | Derive the mathematical expression for ,the angular distance" $\varepsilon$ between the <br> two images seen through the telescope (first as a function of $\theta$ and $\Delta n$, then <br> as a function of $n$ and $\Delta n)$ and calculate its numerical value. | 2.00 p |
| :--- | :--- | :--- |


| A3 | If the focal distance of the telescope's objective is $f_{o b}=0.40 \mathrm{~m}$, derive the <br> linear distance $y$ between the two images, seen in the focal plane of the <br> objective and calculate its numerical value. | 0.50 p |
| :--- | :--- | :--- |

## B. Refraction, but. . .mostly reflection

## B1. Total reflection in geometrical optics

Total reflection occurs when light travels from a medium with refractive index $n_{1}$ to another one, with the refractive index $n_{2}<n_{1}$, at an incidence angle $\theta_{1} \geq l$, where $l$ is the critical value of the incidence angle, called limit angle, beyond which there will be no refracted light. At total reflection, the entire energy of the incident light beam goes to the reflected beam.

| $\mathbf{B 1}$ | Derive the mathematical expression for the limit angle | 0.25 p |
| :--- | :--- | :--- |

## B2. Total reflection in electromagnetic optics

Electromagnetic optics proves that, besides being totally reflected, the incident light beam also penetrates the less refringent medium as an evanescent wave.


Fig. 2
The characteristics of the reflected and the refracted light beams depend on the angle of incidence, as well as on the orientation of the electric field of the light wave (called polarization). For simplicity, let us consider that the electric field intensity is perpendicular on the incidence plane, as represented in Fig. 2. The indices $i, r$, and $t$ refer to the incident, reflected and transmitted properties of light waves while $\vec{k}$ is the wave vector, giving the light propagation orientation. Moreover, $\hat{x}, \hat{y}$ and $\hat{z}$ are the unit vectors of the chosen Cartesian reference frame.
Physical note: The perturbation produced by a plane, monochromatic wave in a point in space at a certain moment of time can be written as $\vec{E}(\vec{r}, t)=\vec{E}_{0} \cos (\omega t-\vec{k} \cdot \vec{r})$, or, to simplify calculations, in the complex form $\vec{e}(\vec{r}, t)=\vec{E}_{0} e^{i(\omega t-\vec{k} \cdot \vec{r})}$, where $i=\sqrt{-1}$, and then taking only the real part of the result.

Mathematical note: For the complex number $z=a \pm i b, a$ is the real part and $b$ is the imaginary
part. It can be written as $z=\underbrace{\sqrt{a^{2}+b^{2}}}_{=|z|}(\underbrace{\frac{a}{\sqrt{a^{2}+b^{2}}}}_{=\cos \varphi} \pm i \underbrace{\frac{b}{\sqrt{a^{2}+b^{2}}}}_{=\sin \varphi})=|z| e^{ \pm i \varphi}$, where $|z|$ is the modulus of the complex number $z$ and $\tan \varphi=b / a$.

## B2.1. Evanescent wave

Knowing that the incident wave is a plane and monochromatic one, characterized by the equation $\vec{e}_{i}(\vec{r}, t)=\vec{E}_{0 i} e^{i\left(o t-\vec{k}_{i} \cdot \vec{r}\right)}$, prove that the mathematical expression for the evanescent wave is $\vec{e}_{t}(\vec{r}, t)=\vec{E}_{0 t} e^{-\alpha z} e^{i \varphi}$ and derive the exact expression for the attenuation coefficient $\alpha$ as a function of the incidence angle $\theta_{1}$, the limit angle $l$, and the wavelength $\lambda$ of the incident wave. Also, derive the exact expression for the phase $\varphi$ of the evanescent wave.

## B2.2. Penetration depth

| B2.2 | Derive the mathematical expression of the distance $\Delta z$ from the interface at <br> which the amplitude of the evanescent wave is $e$ times smaller than at the <br> interface, as a function of the incident wavelength $\lambda$ and calculate its <br> numerical value. The first medium is glass $\left(n_{1}=1.6\right)$, the second is air <br> $\left(n_{2}=1.0\right)$, and the incidence angle of light is $\theta_{1}=40^{\circ}$. | 0.50 p |
| :--- | :--- | :--- |

## B2.3. The phase speed of the evanescent wave

| B2.3 | Derive the mathematical expression for the ratio $\frac{v_{e}}{v_{1}}$, where $v_{e}$ is the phase <br> speed of the evanescent wave and $v_{1}$ - the phase speed of the incident wave, <br> and compute its numerical value for the case of the incidence angle of light <br> of $\theta_{1}=40^{\circ}$. | 0.75 p |
| :--- | :--- | :--- |

## B2.4. The energy transferred from the incident wave to the totally reflected wave

For any value of the incidence angle, the relationship between the amplitude of the field of the reflected wave and that of the incident wave was derived by the French physicist Augustin Fresnel (1788-1829):

$$
E_{0 r}=\frac{n_{1} \cos \theta_{1}-n_{2} \cos \theta_{2}}{n_{1} \cos \theta_{1}+n_{2} \cos \theta_{2}} E_{0 i}
$$

## Romanian Master of Physics 2017

Physical note: If the perturbation produced by a wave in a point in space at a given moment is expressed using complex numbers, then the wave intensity has the mathematical expression $I=\frac{1}{2} \varepsilon_{0} c E_{0}^{*} E_{0}=\frac{1}{2} \varepsilon_{0} c\left|E_{0}\right|^{2}$, where $E_{0}^{*}=a-i b$ is the complex conjugate of the complex number $E_{0}=a+i b$. Here $\varepsilon_{0}$ is the vacuum permittivity and $c$ is the speed of light in vacuum.

B2.4
Prove that the totally reflected wave has the same intensity as the incident wave.
0.50 p

## B2.5. The Goos - Hänchen effect

When an incident wave beam with a finite cross section undergoes total reflection at an interface between two media, the totally reflected wave beam is laterally displaced, on a distance $D$ (see Fig. 3), that was measured for the first time by Goos and Hänchen in 1947. In Fig. 3, the displacement along the surface is $s$, and the Goos - Hänchen shift is the lateral shift $D$ indicated in the diagram. This is the Goos - Hänchen effect. The explanation of this lateral shift is based on the appearance of the evanescent wave at the interface and its propagation parallel to the interface.


Fig. 3

## B2.5.1. The lateral shift

| B2.5.1 | Derive the mathematical expression for the Goos - Hänchen lateral shift $D$, <br> admitting that the phase difference between the totally reflected wave and <br> the incident one is zero at the interface. Consequently, compute the <br> numerical value of the displacement $s$ along the interface as a function of <br> the wavelength $\lambda$ of the incident light, if the first medium is glass <br> $\left(n_{1}=1.6\right)$, the second is air $\left(n_{2}=1.0\right)$, and the incidence angle of light is <br> $\theta_{1}=40^{\circ}$. | 1.00 p |
| :--- | :--- | :--- |

## B2.5.2. Time needed for the total reflection

An alternative explanation of the Goos - Hänchen shift can be given in terms of the time delay associated with the scattering of a radiation pulse at the interface. The incident radiation pulse is not scattered instantaneously by the surface, but reemerges into medium 1 after a time delay $\tau$, during which the pulse propagates parallel to the surface and is displaced by the distance $s$.

| B2.5.2. | Derive the mathematical expression for the time delay $\tau$ and calculate its <br> value if the first medium is glass $\left(n_{1}=1.6\right)$, the second is air $\left(n_{2}=1.0\right)$, <br> the incidence angle of light is $\theta_{1}=40^{\circ}$, and the monochromatic radiation <br> has the wavelength $\lambda=579.1 \mathrm{~nm}$. The light speed in vacuum is <br> $c=3.0 \cdot 10^{8} \mathrm{~m} / \mathrm{s}$. | 0.75 p |
| :--- | :--- | :--- |

## Problem I

## Reflection and refraction of fight

## Answer sheet

## A. $\mathcal{A n}$ interesting prism

|  | Mathematical expression | Numerical value |
| :--- | :---: | :---: |
| A1 | $\theta(n)=\quad \theta(n=1.6)=$ |  |
|  |  |  |


|  |  |  | Numerical value |
| :--- | :--- | :--- | :--- |
|  | $\varepsilon(\theta, \Delta n)=$ | $\varepsilon(n, \Delta n)=$ | $\varepsilon(n, \Delta n)=$ |


|  | Mathematical expression | Numerical value |
| :---: | :---: | :---: |
|  | A3 | $y\left(f_{o b}, n, \Delta n\right)=$ |
|  |  |  |
|  |  |  |

B. Refraction, 6ut. . . mostly reflection

## B1. Total reflection in geometrical optics

$\square$

## Romanian Master of Physics 2017

## B2. Total reflection in electromagnetic optics

## B2.1. Evanescent wave

|  |  |  |
| :--- | :--- | :--- |
| B2.1 | $\alpha\left(\theta_{1}, l, \lambda\right)=\quad \varphi=$ |  |

## B2.2. Penetration depth

|  |  | Numerical value |
| :--- | :--- | :--- |
|  | $\Delta z(\lambda)=\quad \Delta z(\lambda)=$ |  |

## B2.3. The phase speed of the evanescent wave

|  | Mathematical expression | Numerical value |
| :---: | :--- | :--- |
|  |  |  |
| $\mathbf{B 2}^{2} 3$ | $\frac{v_{e}}{v_{1}}=$ | $\frac{v_{e}}{v_{1}}=$ |
|  |  |  |

B2.4. The energy transferred from the incident wave to the totally reflected wave

| B2.4 |  |
| :--- | :--- | :--- |
|  |  |
|  |  |
|  |  |
|  |  |

$\square$
B2.5. The Goos - Hänchen effect

## B2.5.1. The lateral shift

|  | Mathematical expression | Numerical value |
| :---: | :---: | :---: |
| B2.5.1 | $D\left(\lambda, \theta_{1}, l\right)=$ | $s(\lambda)=$ |

B2.5.2. Time needed for the total reflection

|  | Mathematical expression | Numerical value |
| :--- | :--- | :--- |
| B2.5.2 | $\tau\left(\lambda, \theta_{1}, l, c, n_{1}\right)=$ |  |
|  |  | $\tau=$ |

